$Room\text{-}Temperature\ Ferromagnetism\ in\ Mn\text{-}Doped\ Semiconducting\ CdGeP_2$

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The chalcopyrite $CdGeP_2$ doped with Mn have been recently found to exhibit room-temperature ferromagnetism. Isovalent substitution of the Cd site is expected, however, to create m_1 , m_2 , m_3 , in

atom is removed. μ_{α}^{a} is the absolute value of the chemical potential of atom α . Since the formation energies are conventionally defined with respect to the elemental solid(s), we express μ_{α}^{a} as the sum of a component due to the element in its most commonly occurring structure μ_{α}^{s} , and an excess chemical potential μ_{α} , i.e., $\mu_{\alpha}^{a} = \mu_{\alpha}^{s} + \mu_{\alpha}$. Here μ_{α}^{s} for P, Ge, Mn, and Cd are the total energies evaluated for the fully optimized elemental solids in the observed crystal structures [14]. If $\Delta H_f(\text{CdGeP}_2)$ is the formation energy of CdGeP₂, then μ_{Cd} and μ_{Ge} are determined by

$$\mu_{\text{Cd}} + \mu_{\text{Ge}} + 2\mu_P \le \Delta H_f(\text{CdGeP}_2).$$
 (2)

Furthermore, $\mu_{\rm Cd} \leq 0$; $\mu_{\rm Ge} \leq 0$, because otherwise the elemental solids will precipitate. The presence of other intervening binary phases, however, further restricts the values of $\mu_{\rm Cd}$ and $\mu_{\rm Ge}$: One must solve Eq. (2) along with the constraints placed by the formation energies $\Delta H_f({\rm Cd}_3{\rm P}_2)$ and $\Delta H_f({\rm GeP})$ of ${\rm Cd}_3{\rm P}_2$ and ${\rm GeP}$:

$$3\mu_{\mathrm{Cd}} + 2\mu_{P} \le \Delta H_{f}(\mathrm{Cd}_{3}\mathrm{P}_{2}), \tag{3}$$

$$\mu_{\text{Ge}} + \mu_P \le \Delta H_f(\text{GeP}),$$
 (4)

to find the allowed range for μ_{Cd} and μ_{Ge} in CdGeP₂. The electrons ionized upon forming a positively charged defect join the Fermi sea so the formation energy increases by $q\epsilon_f$, where ϵ_f is the fermi energy which varies from 0 eV at the valence band maximum (VBM) of the host material to the band gap of the host. Equations (2)–(4) were solved using the experimental values [15,16] of the formation energies for the binary phases Cd₃P₂ (-1.2 eV) and GeP (-0.3 eV), while a value of -1.5 eV, in the same range as other chalcopyrites [17], was used for CdGeP₂.

The allowed range of chemical potentials μ_{Cd} and μ_{Ge} for CdGeP₂ and the binaries Cd₃P₂ and GeP are given in Fig. 1. There are three distinct chemical potential domains where CdGeP₂ can exist: point : Cd rich, Ge poor;

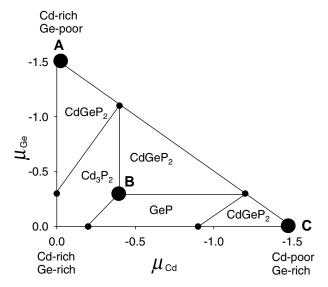


FIG. 1. The range of Cd and Ge chemical potentials where $CdGeP_2$, GeP, and Cd_3P_2 are stable.

point : Cd rich, Ge-rich; and point : Cd poor, Ge rich. Figure 2 shows the formation energies of the intrinsic point defects Ge_{Cd} , V_{Cd} , and V_{Ge} as well as substitutional defects Mn_{Ge} and Mn_{Cd} at the chemical potentials , , and of Fig. 1 as a function of the Fermi energy. The vertical dashed line denotes the generalized gradient approximation (GGA) gap which is underestimated with respect to the experimental 1.72 eV gap. Transition points between charge states are indicated by solid circles. The defects can

in these conditions has no holes and so cannot promote antiferromagnetism. (v)Ge-on-Cd antisite has high formation energy, and would therefore not have appreciable concentration.

Having identified the hole-producing centers that can yield FM, we next examine the predicted solubilities of isolated Mn. Our calculated formation energies for CdGeP₂:Mn and similar calculations for GaAs:Mn show consistently lower values (for the appropriate chemical potentials) in the former case, predicting higher Mn solubility: The lowest formation energy of substituting a Ga atom with Mn in GaAs is 1.0 eV (under Mn-rich, Ga-poor conditions). In contrast, even in the worst-case scenarios, we find a

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