



# Effects of atomic clustering on the optical properties of III-V alloys

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Self-consistent electronic structure calculations together with a structural model are used to study the effect of short-range atomic order on the optical properties of otherwise random  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{As}$ ,  $\text{Ga}_{0.5}\text{In}_{0.5}\text{P}$ , and  $\text{Al}_{0.5}\text{In}_{0.5}\text{As}$  alloys. We find that clustering can reduce the direct band gap of these alloys by as much as 100 meV. Furthermore, sufficiently strong clustering is predicted to transform  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{As}$  into a direct gap material.

$A$  and  $B$  atoms on their fcc sublattice. These deviations take the form of long-range order (LRO), short-range order (SRO), or both. LRO in III-V alloys appears most frequently in the CuPt structure and is accompanied by a reduction in the band gap relative to the disordered phase.<sup>1</sup> This gap reduction reflects zone-folding and level-repulsion<sup>2</sup> and decreases quadratically<sup>3</sup> on the degree  $\alpha$  of LRO. In contrast

Warren-Cowley<sup>4</sup> parameter

$$\alpha_j = 1 - \frac{P_B(j)}{x_B}, \quad (1)$$

random alloy  $P_B(j) = x_B$  and thus  $\alpha_j = 0$  for all atomic shells.

x-ray scattering,<sup>5,6</sup> transmission electron microscopy,<sup>7</sup> and scanning tunneling microscopy.<sup>8</sup> Indirect evidence for SRO comes from nuclear magnetic resonance,<sup>9</sup> resonant Raman scattering,<sup>10</sup> photoluminescence,<sup>11</sup> and photorefraction.<sup>12</sup> Many experiments<sup>5-8,11,13</sup> report clustering-type SRO ( $\alpha > 0$ ), whereas anticlustering seems to be less frequently

tered ( $\alpha > 0$ ) models of  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{As}$ ,  $\text{Ga}_{0.5}\text{In}_{0.5}\text{P}$ , and  $\text{Al}_{0.5}\text{In}_{0.5}\text{As}$  alloys. We find that local clustering can (i) reduce the band gap of III-V alloys to a similar extent as LRO does, (ii) transform the indirect-gap material  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{As}$  into a direct-gap one, and (iii) localizes the band edge wave functions preferentially on a particular type of cluster, which thus acts as a "local quantum-well." We will discuss the chemical trends for wave-function localization in the sequence Al/Ga/In.

The standard approach to the electronic structure of alloys—the virtual crystal approximation—does not distinguish  $A$  from  $B$  atoms irrespective of their chemical disparity. The single-site coherent potential approximation assumes

environments. A more realistic description would allow for a *distribution* of  $A$  and  $B$  atoms reflecting the existence of many distinct local atomic environments in the alloy. The most direct approach to this description would involve application of band theory to fictitious solids with huge ( $\geq 1000$  atom) supercells<sup>16-18</sup> whose sites are occupied by  $A$  and  $B$  atoms according to a prescribed degree (one or finite) of

cal tight-binding<sup>17</sup> and empirical pseudopotentials in a plane-wave basis.<sup>18</sup> These studies showed that the results can be mimicked very well using much smaller unit cells ( $\sim 10$  atoms) but with *specially selected site occupations and cell geometries* (special quasirandom structures, or SQS).

SQSs in the context of self-consistent nonlocal pseudopotential

(LDA) to obtain the electronic energy bands of III-V alloys

type in which the nearest-neighbor shell has  $\alpha_1 > 0$  and all subsequent shells are random. We use 16-atom SQSs with the same computational parameters as described in Ref. 15. The atomic positions of  $\text{Ga}_{0.5}\text{In}_{0.5}\text{P}$  and  $\text{Al}_{0.5}\text{In}_{0.5}\text{As}$  in the

field,<sup>20</sup> while maintaining cubic symmetry for the cell-external degrees of freedom. We are interested in calculating alloy band gaps of a particular zinc-blende representation

of a few SQS states  $|j\mathbf{K}\rangle$ , weighted with the spectral density

and  $j$  are band indices).

$$\langle E_n(\mathbf{k}) \rangle_{\text{alloy}} = \frac{1}{W} \sum_{j, \mathbf{K}} |\langle j\mathbf{K} | n\mathbf{k} \rangle|^2 E_j(\mathbf{K}), \quad (2)$$

where the sum runs over a few SQS states around a peak in the spectral density of the state  $|n\mathbf{k}\rangle$ ,  $\mathbf{K}$  must differ from  $\mathbf{k}$  by a reciprocal lattice vector of the SQS, and  $W$  is a normalization constant.

Table I summarizes our calculated [Eq. (2)] alloy band gaps<sup>21</sup> with and without SRO. We choose in this study a relatively large and positive SRO parameter of  $\alpha = \frac{1}{6}$  in order to emphasize the effect of local clustering. (Note, however,

TABLE I. LDA corrected band gaps in eV (measured from the top of the

randomness ( $\alpha=\eta=0$ ), clustering-type SRO ( $\alpha=\frac{1}{6}$ ), and CuPt-type LRO ( $\eta=1$ ). Here,  $\bar{E}_g$  denotes the average gap of the binaries at their equilibrium volumes. Chemical symbols in parentheses denote the sublattice on which the VBM and CBM are localized, respectively, and  $D$  denotes that the state

lar type of cation tetrahedron  $A_nB_{4-n}$ . We define a cluster

$$w_n(j\bar{\mathbf{K}}) = \frac{1}{N} \sum |\langle \chi_0(\mathbf{r}_i) | j\bar{\mathbf{K}} \rangle|^2, \quad (3)$$

where  $N_n$  is the number of clusters of type  $A_nB_{4-n}$ , the sum runs over all positions of  $A$ ,  $B$ , and  $C$  atoms participating in

System	$\bar{E}_g$	Random	SRO		LRO	
			$\alpha=\frac{1}{6}$		$\eta=1$	
	2.18	2.17	2.16		2.10	

that SRO parameters as large as the present one have been reported in the literature<sup>5</sup>.) The table also gives the average band gaps of the binaries  $\bar{E}_g = \frac{1}{2}[E_g(AC) + E_g(BC)]$  and the calculated gaps of CuPt ordered alloys ( $\eta=1$ ). The optical bowing parameter  $b$  is defined by

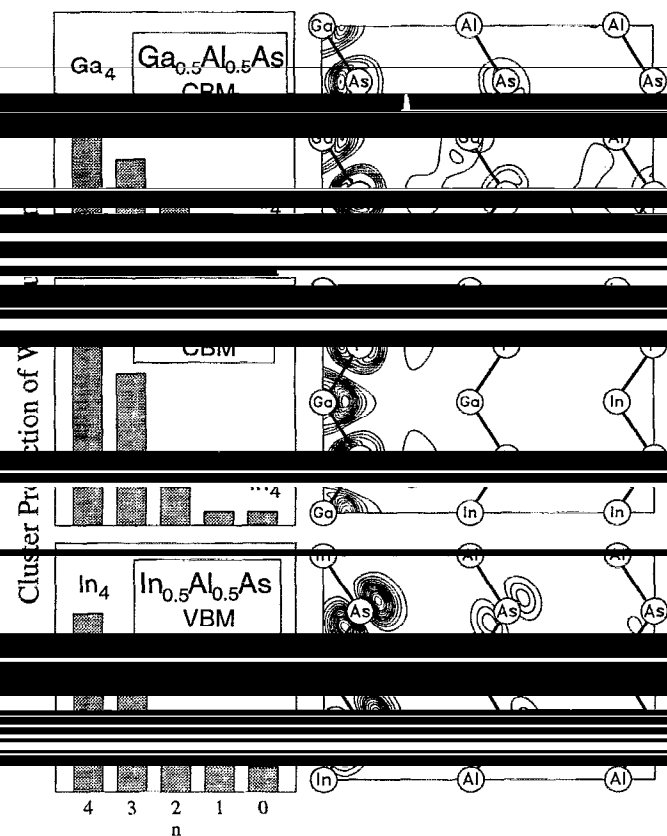
relative to  $E_g$  by 0.05, 0.10, and 0.15 eV for  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{As}$ ,  $\text{Ga}_{0.5}\text{In}_{0.5}\text{P}$ , and  $\text{Al}_{0.5}\text{In}_{0.5}\text{As}$  respectively. This gives bowing parameters of 0.20, 0.40, and 0.52 eV, compared with the experimental results of 0.37 (Ref. 22), 0.70 (Ref. 23), and 0.74 eV (Ref. 24), respectively. (ii) LRO of the CuPt type ( $\eta=1$ ) reduces the band gaps relative to  $\bar{E}_g$  by 0.41, 0.50,

parameter  $\alpha=\frac{1}{6}$  also reduces significantly the direct band

and 1/0 inside the sphere. Histograms of the distribution of  $\text{Ga}_{0.5}\text{In}_{0.5}\text{P}$ , and the valence band maximum of  $\text{Al}_{0.5}\text{In}_{0.5}\text{As}$  are also shown in Fig. 1, demonstrating localization on particular "pure" clusters. (The CBM in  $\text{Al}_{0.5}\text{In}_{0.5}\text{As}$  is weakly localized on In-rich clusters.)

The results of Fig. 1 imply that the clusters act as "impurity-like traps" for electrons or holes, very much like isoelectronic impurities which are known to bind carriers if the difference between their local potential and that of the

ters in the respective alloys spatially extends the range of the



than the  $\kappa$ -like indirect transition. The latter seems to be direct band gap crossover in  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{As}$  with sufficient local clustering. This is in marked contrast to the findings of Fu *et al.* who find *negative* optical bowing upon clustering.<sup>14</sup> Maeri *et al.* obtain the same sign of band gap modifications

parameter assumed there.<sup>15</sup>

In order to understand the physical mechanism leading to the large gap reduction reported in Table I, we analyze the wave functions of the lowest conduction (CBM) and highest valence (VBM) states in the clusters and then in the alloys

either the lowest conduction band edge ( $\text{Al}_{0.5}\text{Ga}_{0.5}\text{As}$ ) or the strongly localize on clusters of one type of cations, while the corresponding states in the random alloys show no such tendency. This is illustrated by plotting contours of the wave

segregation (i.e., localization) can be quantified by projecting the wave functions on atomic spheres belonging to a particu-

we show contour plots of the corresponding wave function squared on the (110) plane. The states are clearly segregated on one type of cluster

perturbing potential, thus enhancing its strength and leading to binding. Similarly, an isolated In impurity in GaP or AIAs will probably not bind a hole, but the corresponding clusters do. For sufficiently large cluster sizes we can rephrase the preceding argument in terms of band theory: The segregation of the band edge wave functions is dictated by the band

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- <sup>4</sup>J. M. Cowley, *J. Appl. Phys.* **21**, 24 (1950).
- <sup>5</sup>H.-G. Brühl, L. Hildisch, H. Morwinski, W. Schmidt, and E. Schubert,

Common anion volume  $V$  reverses the order of the lowest

only  $\sim 0.25$  eV higher<sup>27</sup> than the one in expanded GaP.) Thus, the lowest conduction wavefunction in a clustered  $\text{Ga}_{0.5}\text{In}_{0.5}\text{P}$  alloy should segregate on the Ga-rich clusters as found in the calculation (Fig. 1). For electrons, the observed

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- <sup>15</sup>R. Magri, S. Froyen, and A. Zunger, *Phys. Rev. B* **44**, 7947 (1991).

9.16, 8.46, and 7.85 eV, respectively). The hole localization, however, is less directly correlated with an atomic property of the cations alone, since the wavefunctions of the

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in common-anion alloys can reduce the energy gap considerably with respect to the ideal random alloys. This is accompanied by localization of the band edge wave functions

for each binary constituent. For the alloys we use the average of these shifts, which are 0.84, 0.61, and 0.77 eV, for  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{As}$ ,  $\text{Ga}_{0.5}\text{In}_{0.5}\text{P}$ , and  $\text{Al}_{0.5}\text{In}_{0.5}\text{As}$ , respectively. Note that energy differences such as optical bowing coefficients are not affected by this procedure.

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