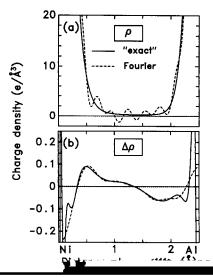
with the results of the self-consistent *ab initio* calculations based on the local density formalism [12, 13]. We focus on the following questions:

(i) How well can *ab initio* band theory describe the first few (low-angle) structure

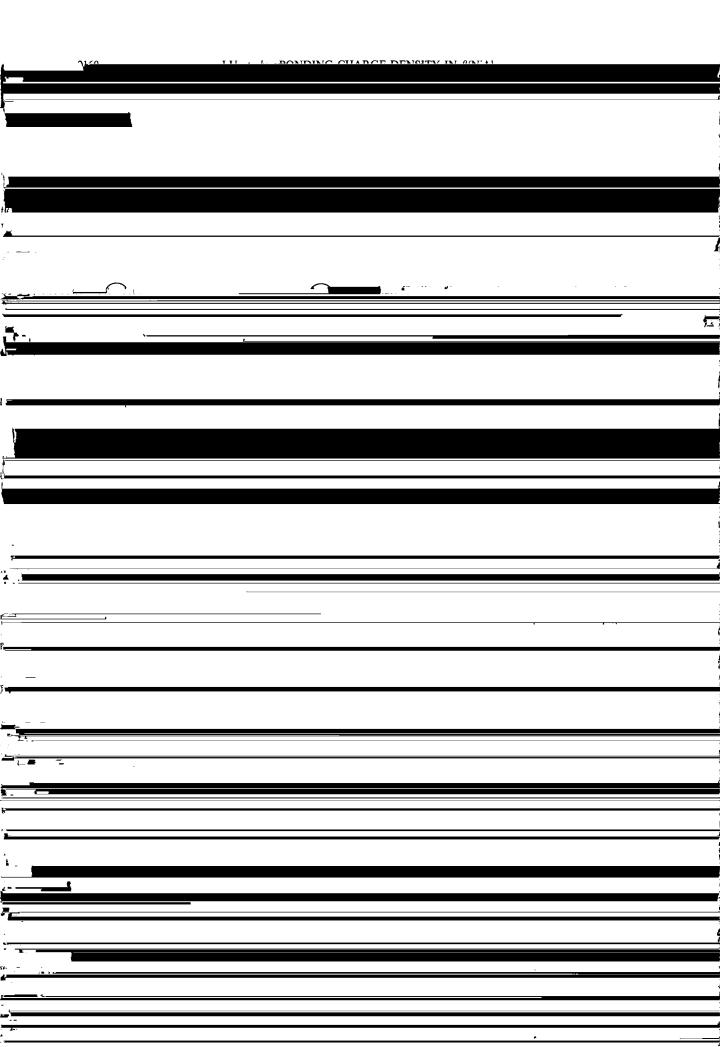
where the result naturally depends on the highest momentum (G_{max}) included in this sum (as we will see below, current high-precision experiments are limited to rather small cut-off values G_{max}). If the temperature factor can be deconvoluted from



representation the *total* density ρ [Fig. 1(a)] still exhibits significant oscillations in the bonding region. On the other hand, the density difference $\Delta \rho_{\rm sup}({\bf r},G_{\rm big})$ (equation (12) and Fig. 1(b)] obtained from a Fourier series using $G_{\rm big}=2\pi/a(6,3,1)$ closely mimics the directly calculated $\Delta \rho_{\rm sup}({\bf r})$ in the bonding regions [of course, $\Delta \rho_{\rm sup}({\bf r},G_{\rm big})$ still fails to reproduce the nodal structure near the core]. Note that the maximum magnitude of the static deformation density $\Delta \rho_{\rm sup}({\bf r})$ outside the core is only $\sim 0.1~{\rm e/\AA^3}$, while the total density $\rho({\bf r})$ has a magnitude of $\sim 10~{\rm e/\AA^3}$ at this point. Clearly, the bonding charge is tiny.

Figure 2 shows as solid line the calculated static density deformation $\Delta \rho_{\text{sup}}(\mathbf{r})$ calculated without any Fourier truncation, comparing it to $\Delta \rho_{\text{sup}}(\mathbf{r}, G_{\text{big}})$ of equation (12), in which a large but finite cut-off $G_{\text{constant}} = 2\pi J a(6, 3, 1) (54 \text{ stars})$ was used. We see a sain

that while the Fourier representation rounds off the



superposition atomic structure factors (using

Here, $B_{\rm Ni}$ and $B_{\rm Al}$ are adjustable parameters while $\rho_{\rm I}$ and $\rho^{\rm MT}$ are fixed by theory (Table 1). This yields rather reasonable values of $B_{\rm Ni}=0.56$ and $B_{\rm Al}=0.71$. Hartree-Fock data [17] as an example) are $F_{\text{sup}}(100) = 13.30$ and $F_{\text{sup}}(200) = 22.84$. Hence the 11

<u>.</u>	accurate measurements of crystalline structure factors to date. We also include in this table the "forbidden" (222) reflection measured by Alkire et al.	Static $\Delta \rho(\mathbf{r})$ Dynamic $\Delta F(\mathbf{r})$ $(100)+(110)+(111)+(200)$ $(100)+(110)+(111)+(200)$ $(100)+(110)+(111)+(200)$ $(100)+(110)+(111)+(200)$ $(100)+(110)+(111)+(200)$ $(100)+(110)+(111)+(200)$	
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